Frame bare

GCG as a scalar

field model

Hamilton-Jacobi Formalism for Chaplygin Gasses

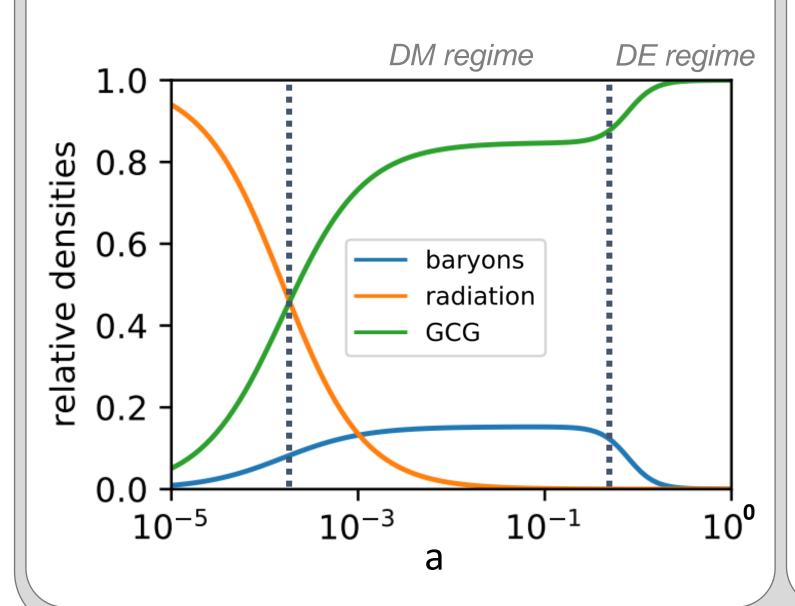
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1. Background

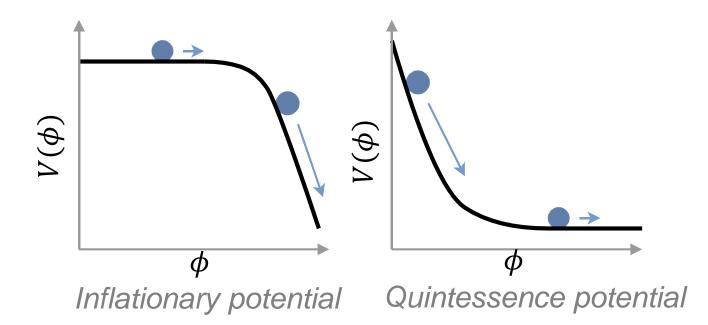
A Generalised Chaplygin Gas (GCG) is a cosmological fluid that seeks to unify Dark Matter (DM) and Dark Energy (DE) [1]. DE domination is characterised by the exponential expansion of the universe.

When the universe is small the GCG behaves as DM, switching its behaviour to DE as it grows.



Scalar fields are commonly used to model inflation and (more recently) quintessence (DM and/or DE). When the field is moving slowly in its potential, it causes the universe to undergo accelerated expansion.

The field obeys a similar Equation Of Motion (EOM) as a ball on a frictional slope. When rolling slowly, the field acts like the *inflaton/DE*.



The Hamilton-Jacobi (HJ) formalism [2] uses the field's value ϕ as the time parameter. This reduces the problem from one second order EOM to two first order equations.

Solve the bare scalar model in the HJ formalism Add a nonzero potential possibly radiation) Add baryonic matter (and possibly radiation)

Scalar field Lagrangians with **non-canonical kinetic terms** are motivated by some beyond-Standard Model theories. One of these replicates the behaviour of the standard GCG [3]:

$$\mathcal{L}(\dot{\phi}) = -A^{\frac{1}{1+\alpha}} \left(1 - |\dot{\phi}|^{\frac{1+\alpha}{\alpha}} \right)^{\frac{\alpha}{1+\alpha}}$$

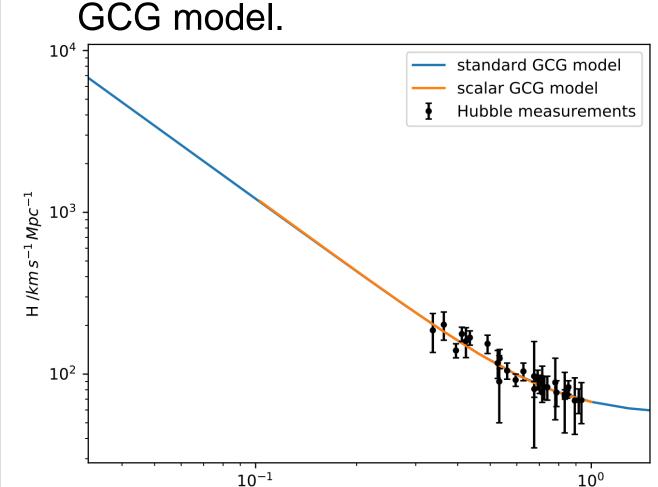
We show this equivalence and derive the **Friedmann equations** for a non-canonical scalar field model in the HJ formalism, following Binétruy et. al. [4].

The model is generalised further by adding a potential term to the Lagrangian, and suitable parameterisations are explored.

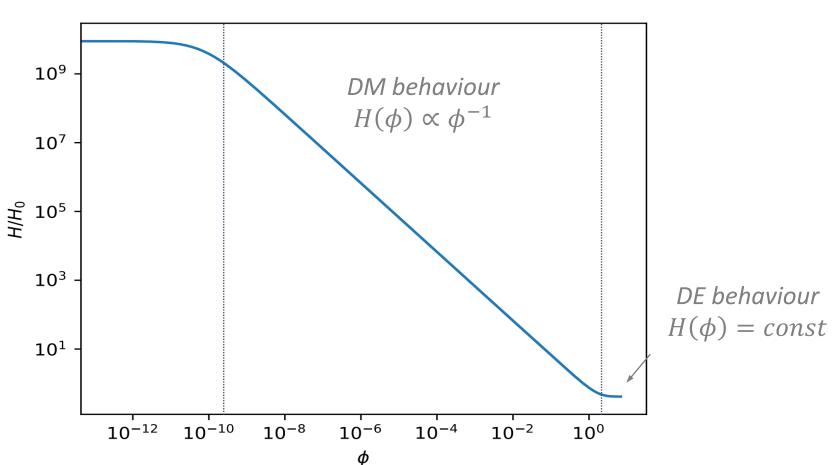
Matter and radiation are added, and the validity of the parameterisations is re-evaluated.

3. Work Undertaken

(a). Framed scalar GCG model in the HJ formalism, solving it numerically. Agrees perfectly with the standard GCG model



(b). Added a potential, parameterising it indirectly by parameterising the Hubble parameter, $H(\phi)$.



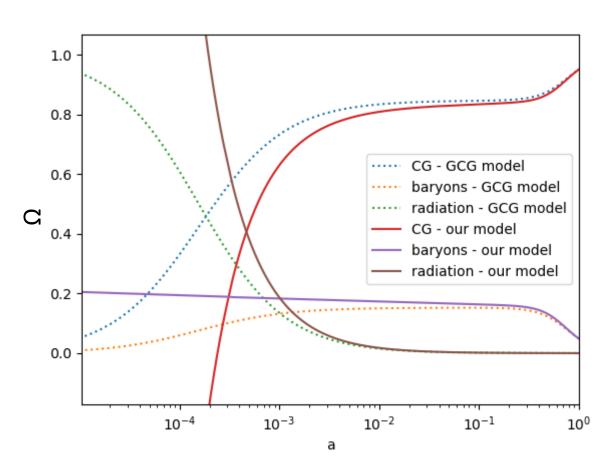
(c). Varied the parameters of the potential to visualise their effects $H(\phi)$ in t

on the model. $p_3 = 2.25$ $p_3 = 2.45$ $p_3 = 2.65$ --- $p_3=2.75$ Hubble data 10^{2} 2.0 $p_3 = 2.25$ $p_3 = 2.35$ $p_3 = 2.45$ 1.5 $p_3 = 2.55$ -- $p_3=2.65$ --- $p_3=2.75$ 0.5 10^{-1}

We based the parameterisation on behaviour of $H(\phi)$ in the **zero-potential case** (above):

$$H(\phi) = \frac{p_1}{\phi} \left(1 + \left(\frac{\phi}{p_2} \right)^{p_3} \right)^{\frac{1}{p_3}}$$

(d). Added baryonic matter & radiation to the model.



Fits standard GCG model (dotted lines) well for DE and DM. Diverges as expected in radiation regime – parameterisation is not valid there.

References

- [1] A. Y. Kamenshchik, U. Moschella and V. Pasquier, An Alternative to quintessence, Phys. Lett. B511(2001) 265 [gr-qc/0103004].
- [2] D. S. Salopek and J. R. Bond, Nonlinear evolution of long wavelength metric fluctuations in inflationary models, Phys. Rev. D42(1990) 3936. [3] M. Bento, O. Bertolami and A. Sen, Generalized Chaplygin gas, accelerated expansion and dark energy matter unification, Phys. Rev. D66(2002) 043507 [gr-gc/0202064]
- [4] P. Binétruy, J. Mabillard and M. Pieroni, Universality in generalized models of inflation, Journal of Cosmology and Astroparticle Physics 2017(2017) 060