

How Universal can Turbulence be ?



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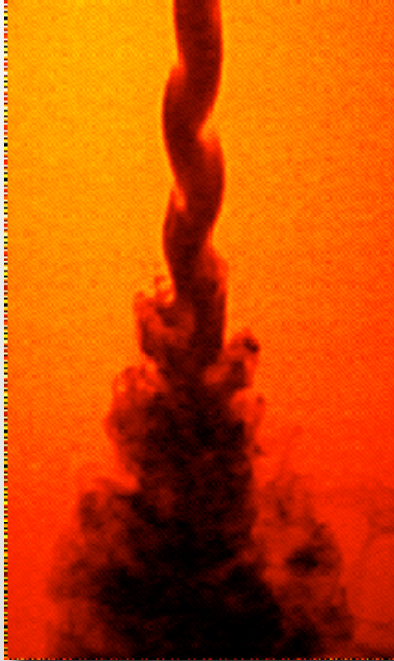
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content

- **comments on** universality
- scaling properties
- characterization without scaling assumption
- how universal is turbulence?

How Universal can Turbulence be ?



Unsolved problems in physics

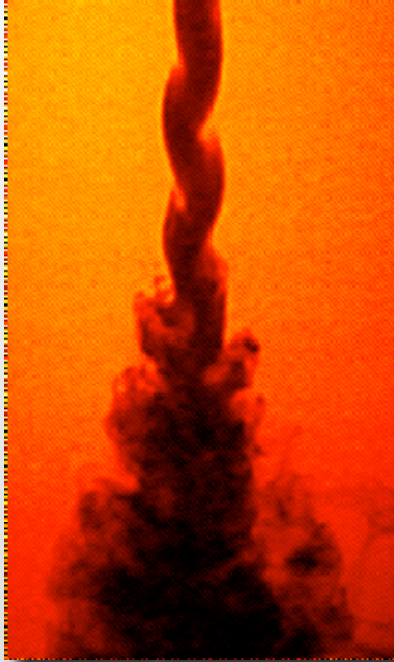
Is it possible to make a theoretical model to describe the behavior of a turbulent flow — in particular, its internal structures?



concept of universality

- only one model is required

concept of universality



- Kolmogorov 41
inertial range

$$\eta \ll r \ll L$$

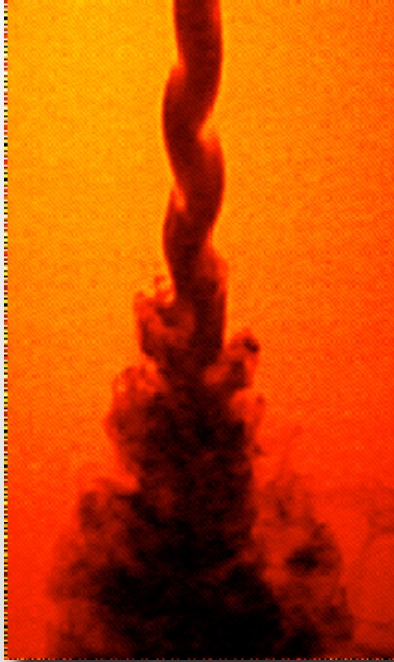
- Kolmogorov Obukhov 62
self-similar energy statistics

$$\langle (\ln \epsilon_r)^2 \rangle = \Lambda_0^2 - \mu \ln(r)$$

scaling law

$$\begin{aligned} \langle u_r^n \rangle &\propto \langle \epsilon_r^{n/3} \rangle r^{n/3} \\ &\propto r^{\xi_n} \end{aligned}$$

concept of universality



- Kolmogorov 41
inertial range

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scaling law ??

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content

- comments on
 - universality is based on what?
- scaling properties

scaling properties

hints for scaling behaviour

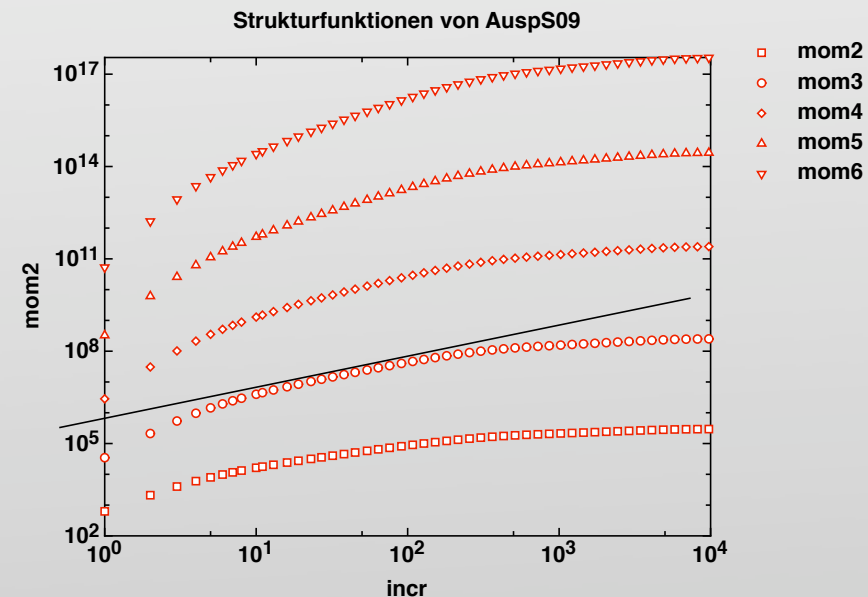
- 4/5th law
$$\langle u_r^3 \rangle = -\frac{4}{5} \langle \epsilon_r \rangle r + 6\nu \frac{\partial}{\partial r} \langle u_r^2 \rangle$$

scaling properties

hints for scaling behaviour

- 4/5th law
- structure functions

$$\langle u_r^n \rangle \propto r^{\xi_n}$$

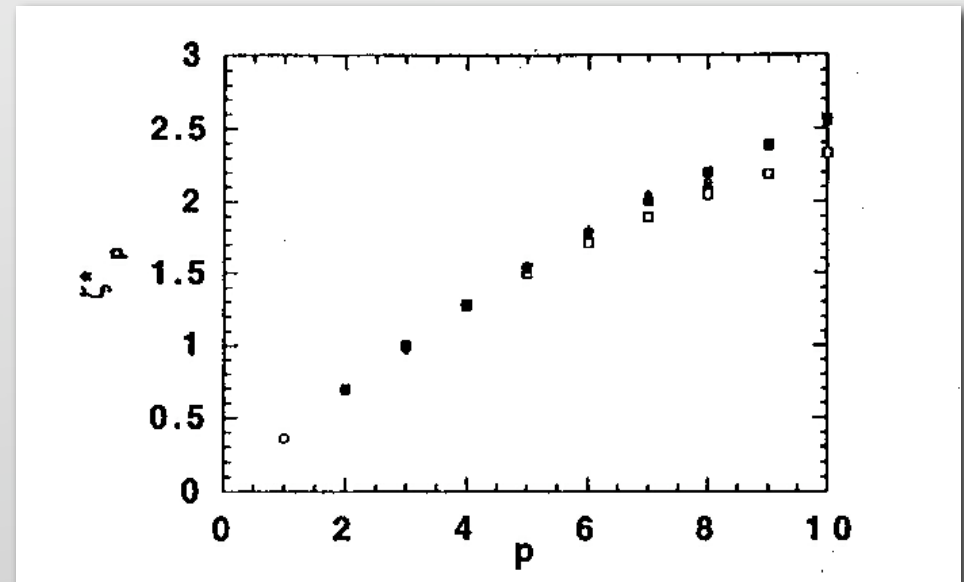
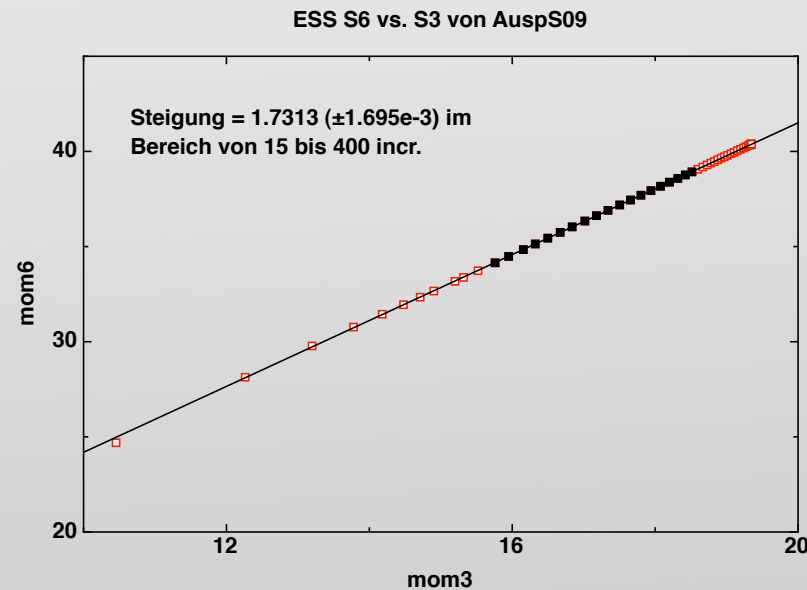


scaling properties

hints for scaling behaviour

- 4/5th law
- structure functions
- ESS

$$\langle u_r^n \rangle \propto \langle u_r^3 \rangle^{\xi_n}$$



(Benzi et.al. (1993))

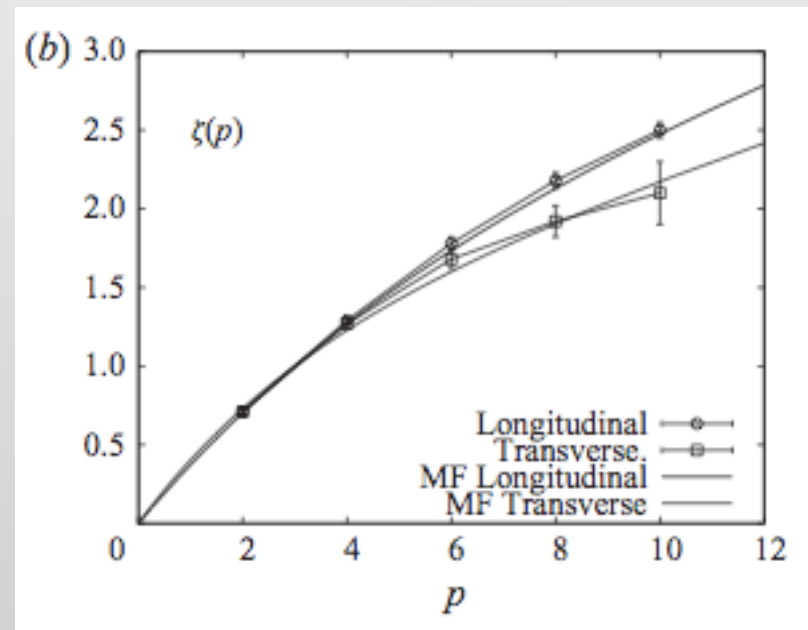
A.Arneodo et.al. Europhys. Lett (1996)

scaling properties

hints for scaling behaviour

- 4/5th law
- structure functions
- ESS
- long / transversal

$$\langle u_r^n \rangle \propto r^{\xi_n}$$



Benzi et.al. JFM (2010)

scaling properties

are there hints **against** scaling behaviour ?

- Castaing

$$\langle (\ln \epsilon_r)^2 \rangle \propto r^\beta$$

Castaing et.al. PRL (1994)

- Kolmogorov Obukhov 62
self-similar energy statistics

$$\langle (\ln \epsilon_r)^2 \rangle = \Lambda_0^2 - \mu \ln(r)$$

scaling properties

are there hints against scaling behaviour ?

- Castaing
- Re dependencies
- ...

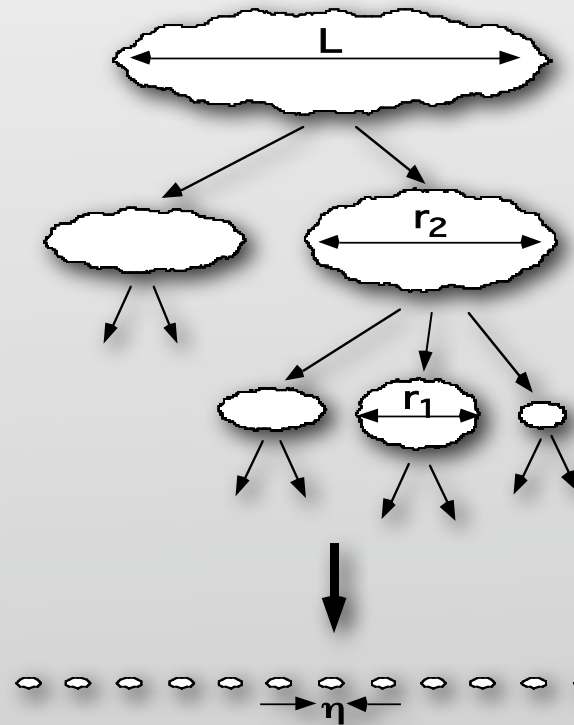
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stochastic cascade process

○ idea of a turbulent cascade:

large vortices are generating small ones



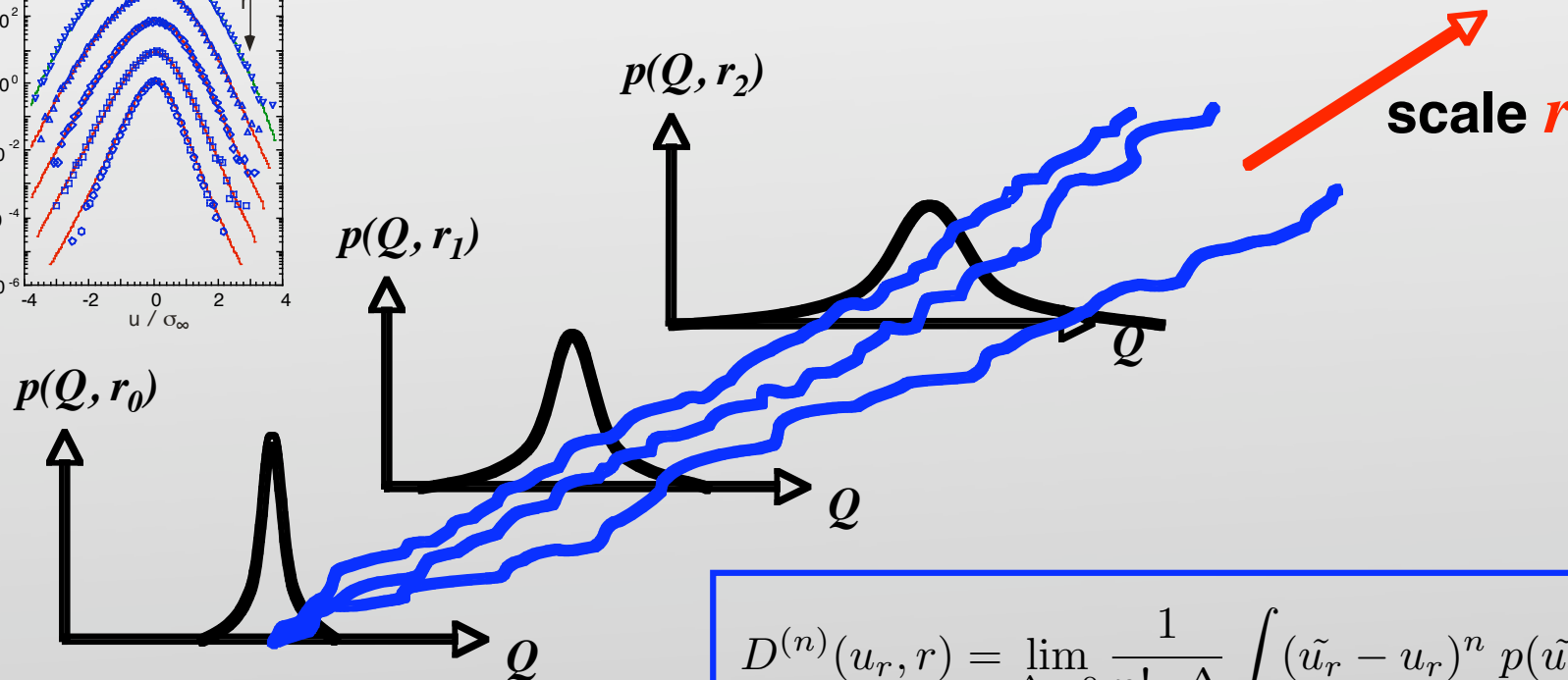
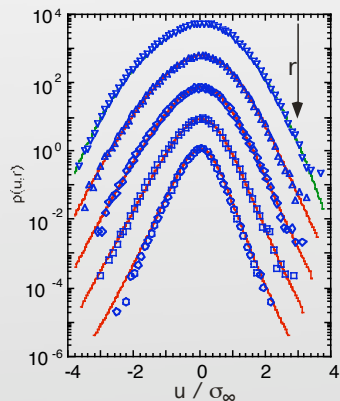
$$\partial_r u_r$$

$$\partial_r p_r(u_r)$$

=> stochastic cascade process evolving in r

stochastic cascade the idea

- the reconstruction of stochastic equation goes back to Kolmogorov work from 1931



$$D^{(n)}(u_r, r) = \lim_{\Delta \rightarrow 0} \frac{1}{n! \cdot \Delta} \int (\tilde{u}_r - u_r)^n p(\tilde{u}_r, r + \Delta | u_r, r) d\tilde{u}_r$$

$$-\frac{\partial}{\partial r} p(u_r, r | u_0, r_0) = \left[-\frac{\partial}{\partial u_r} D^{(1)}(u_r, r) + \frac{\partial^2}{\partial u_r^2} D^{(2)}(u_r, r) \right] p(u_r, r | u_0, r_0)$$

measured Fokker-Planck equation further results

$$-\frac{\partial}{\partial r}p(u_r, r|u_0 r_0) = \left[-\frac{\partial}{\partial u_r}D^{(1)}(u_r, r) + \frac{\partial^2}{\partial u_r^2}D^{(2)}(u_r, r)\right]p(u_r, r|u_0, r_0)$$

differential equations for the structure function

$$-\frac{\partial}{\partial r}\langle u_r^n \rangle = n \cdot \langle u_r^{n-1} D^{(1)}(u_r, r) \rangle + n \cdot (n-1) \langle u_r^{n-2} D^{(2)}(u_r, r) \rangle$$

measured Fokker-Planck equation further results

$$-\frac{\partial}{\partial r} p(u_r, r | u_0 r_0) = \left[-\frac{\partial}{\partial u_r} D^{(1)}(u_r, r) + \frac{\partial^2}{\partial u_r^2} D^{(2)}(u_r, r) \right] p(u_r, r | u_0, r_0)$$

differential equations for the structure function

$$-\frac{\partial}{\partial r} \langle u_r^n \rangle = n \cdot \langle u_r^{n-1} D^{(1)}(u_r, r) \rangle + n \cdot (n-1) \langle u_r^{n-2} D^{(2)}(u_r, r) \rangle$$

- closed equation for structure functions and scaling behaviour if

$$D^{(1)}(u_r, r) = d_1^u(r) u_r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r) u_r + d_2^{uu}(r) u_r^2$$

- and $d_1^u(r) \propto 1/r$ $d_2^{uu}(r) \propto 1/r$ other terms zero => K62

differential equations for the structure function

$$-\frac{\partial}{\partial r} \langle u_r^n \rangle = n \cdot \langle u_r^{n-1} D^{(1)}(u_r, r) \rangle + n \cdot (n-1) \langle u_r^{n-2} D^{(2)}(u_r, r) \rangle$$

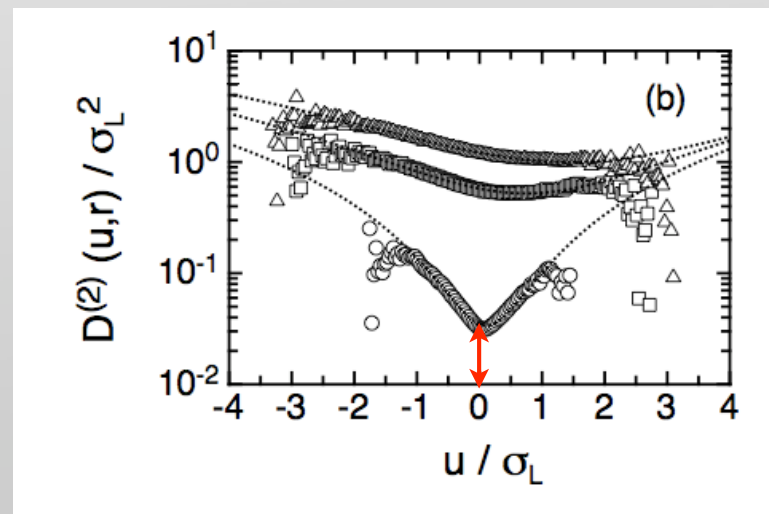
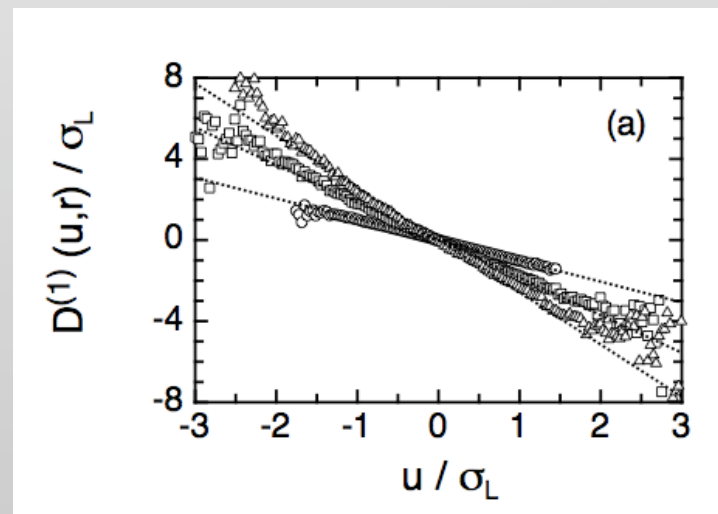
- violation of proper scaling

$$D_1^{(1)}(u_r, r) = d_1^u(r) u_r$$

$$d_1^u(r) \propto 1/r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r) u_r + d_2^{uu}(r) u_r^2$$

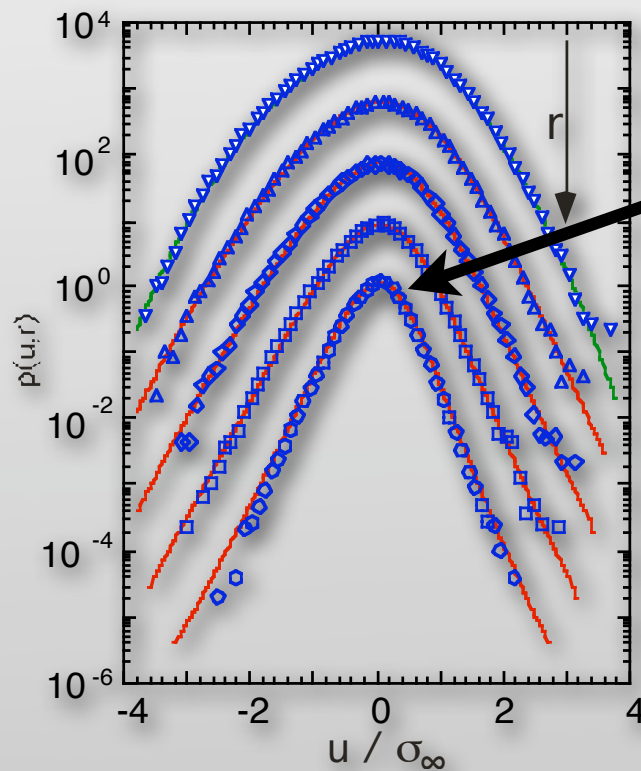
$$d_2^{uu}(r) \propto 1/r$$



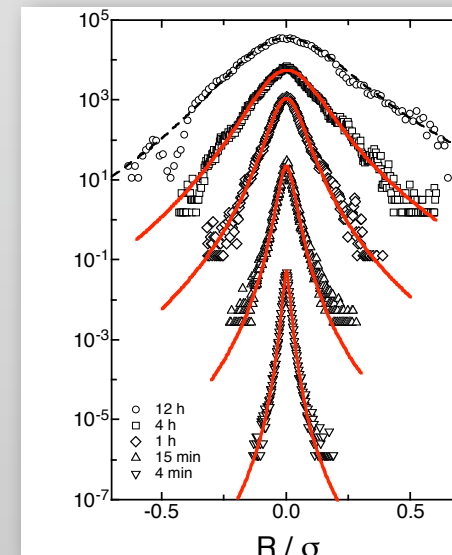
additive term in the diffusion term: -> additive noise

$$D_1^{(1)}(u_r, r) = d_1^u(r)u_r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r)u_r + d_2^{uu}(r)u_r^2$$



Gaussian tip



comparison of data with
numerical solution of the
Kolmogorov equation

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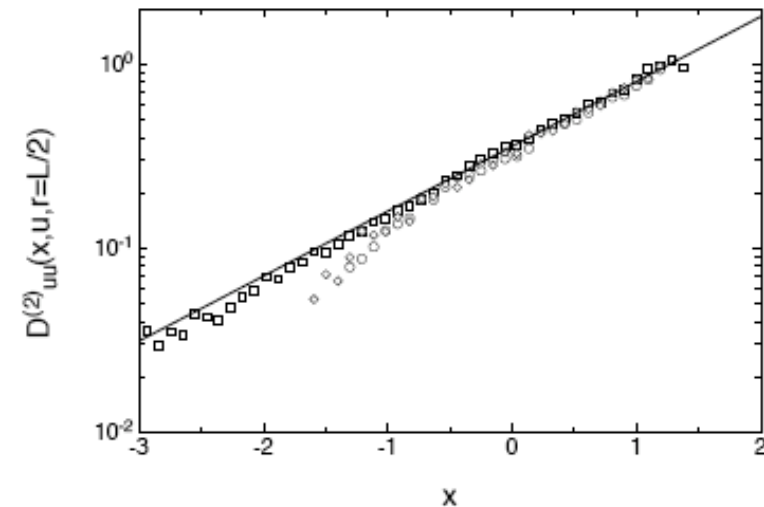
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 - energy cascade

2-dim Fokker-Planck analysis

$$\mathbf{q}(r) = \begin{pmatrix} u(r) \\ x(r) \end{pmatrix}.$$

$$x(r) = \ln(\epsilon_r/\bar{\epsilon}).$$



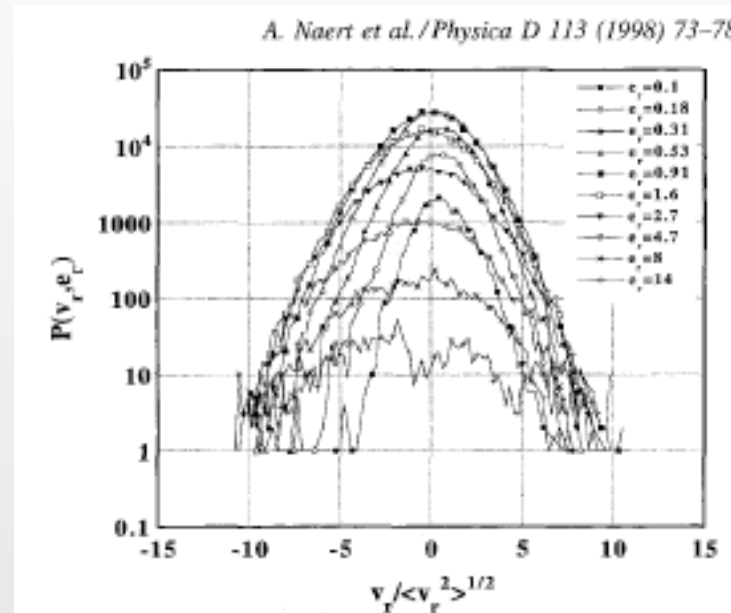
$u = -\sigma_\infty$ (circles), $u = 0$ (squares) and $u = +\sigma_\infty$ (diamonds).

result in form of a
Langevin equation for the cascade

$$\begin{aligned} -\frac{\partial}{\partial r} u(r) &= -\frac{1}{r} \gamma(r) u(r) + m \exp\left(\frac{a_1}{2} x(r)\right) \Gamma_u(r), \\ -\frac{\partial}{\partial r} x(r) &= +\frac{1}{r} G(r) x(r) + \frac{1}{r} F(r) + \sqrt{\frac{1}{r} D_{xx}^{(2)}(u, x, r)} \Gamma_x(r). \end{aligned}$$

no multiplicative
noise anymore !!
No intermittency!!

energy cascade -2-



result in form of a
Langevin equation for the cascade

$$\begin{aligned} -\frac{\partial}{\partial r} u(r) &= -\frac{1}{r} \gamma(r) u(r) + m \exp\left(\frac{a_1}{2} x(r)\right) \Gamma_u(r), \\ -\frac{\partial}{\partial r} x(r) &= +\frac{1}{r} G(r) x(r) + \frac{1}{r} F(r) + \sqrt{\frac{1}{r} D_{xx}^{(2)}(u, x, r)} \Gamma_x(r). \end{aligned}$$

scaling violating term - necessary for
energy conservation

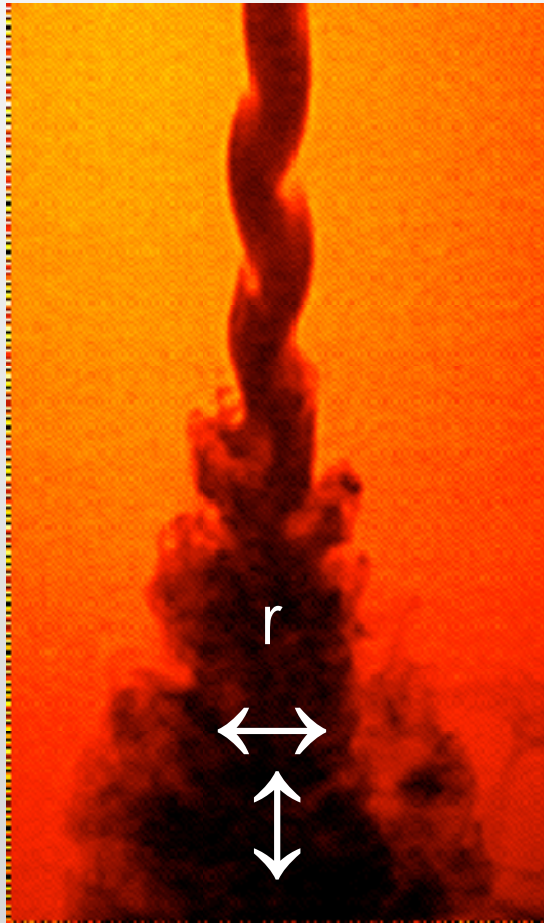
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- scaling properties
- **characterization without scaling assumption**
 - multi-scale statistics - **additive noise**
 - energy cascade and **energy conservation**

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turbulence: long/transversal



spatial correlation in different directions

Quantities

- **longitudinal** increment

$$u_r(x) = [\vec{u}(\vec{x} + \vec{r}) - \vec{u}(\vec{x})] \cdot \hat{r}$$

- **transversal** increment

$$v_r(x) = |[\vec{u}(\vec{x} + \vec{r}) - \vec{u}(\vec{x})] \times \hat{r}|$$

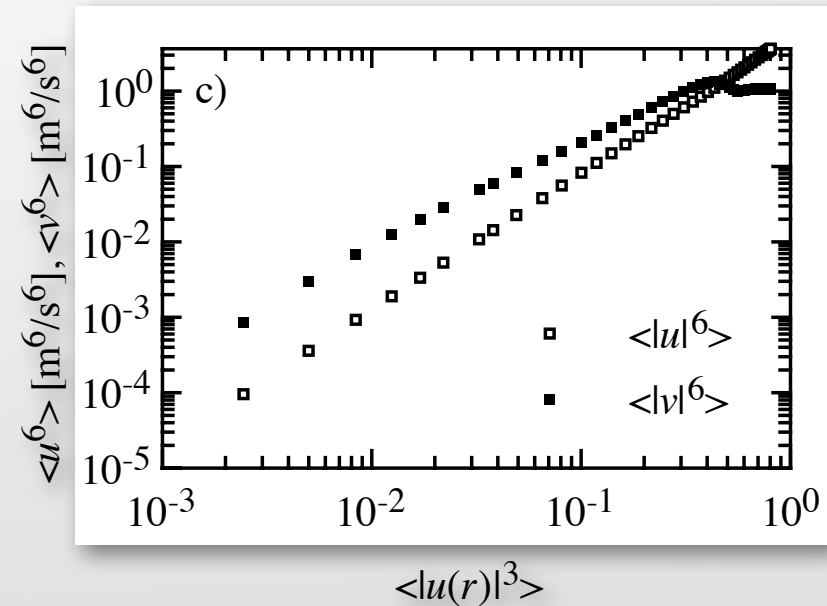
turbulence: long/transversal -2-

extended selfsimilarity **ESS**

supposed scaling laws

$$\langle |u_r|^n \rangle \propto \langle |u_r|^3 \rangle^{\xi_n^l}$$

$$\langle |v_r|^n \rangle \propto \langle |u_r|^3 \rangle^{\xi_n^t}$$



turbulence: long/transversal -2-

extended selfsimilarity **ESS**

supposed scaling laws

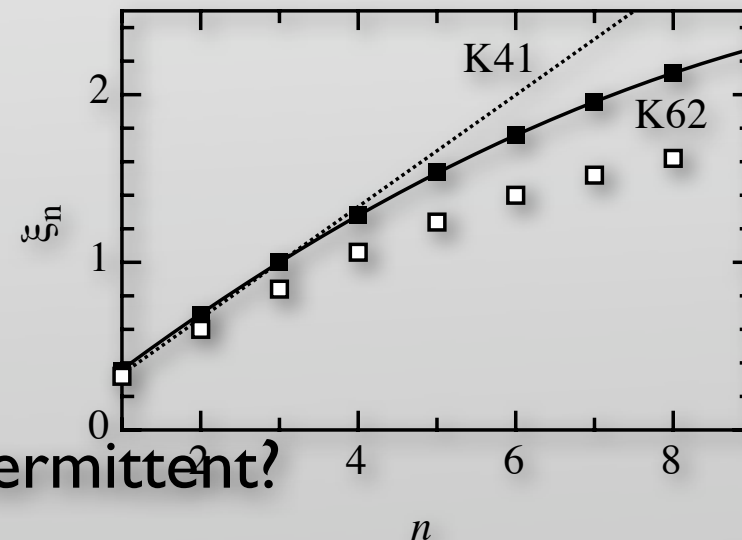
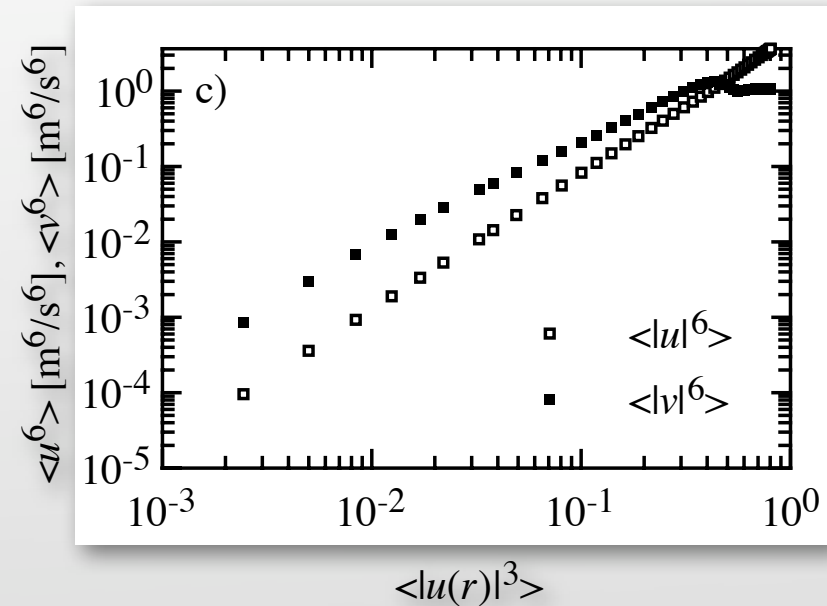
$$\langle |u_r|^n \rangle \propto \langle |u_r|^3 \rangle^{\xi_n^l}$$

$$\langle |v_r|^n \rangle \propto \langle |u_r|^3 \rangle^{\xi_n^t}$$

open problem: (Antonia 97, Benzi 97,
van der Water 99, Grossman et.al. 97....)

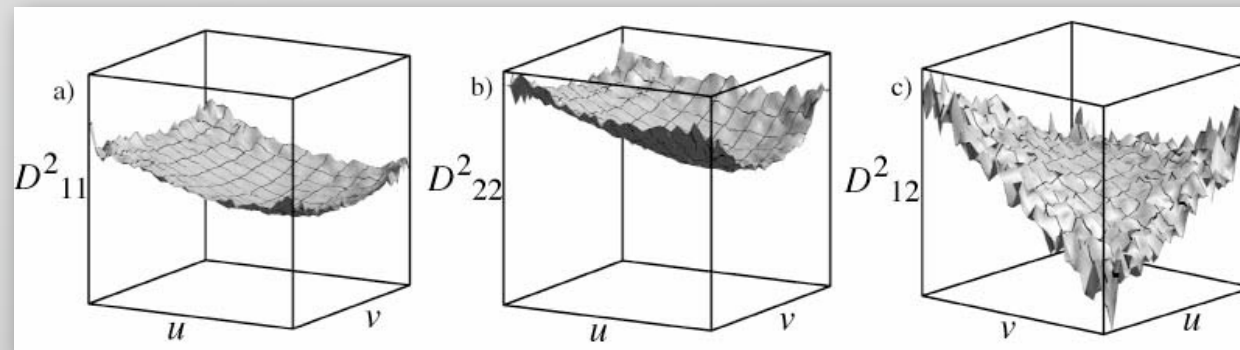
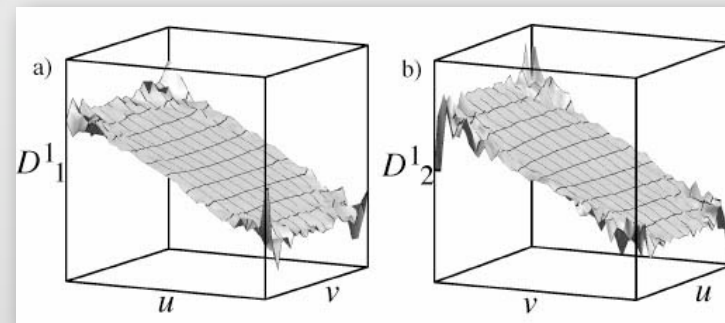
$$\xi_n^t > \xi_n^l$$

are transversal structures more intermittent?



turbulence: long/transversal -3-

$$-r \frac{\partial}{\partial r} p(\mathbf{u}, r | \mathbf{u}_0, r_0) = \left(- \sum_{i=1}^n \frac{\partial}{\partial u_i} D_i^{(1)} + \sum_{i,j=1}^n \frac{\partial^2}{\partial u_i \partial u_j} D_{ij}^{(2)} \right) p(\mathbf{u}, r | \mathbf{u}_0, r_0)$$



turbulence: long/transversal -3-


$$-r \frac{\partial}{\partial r} p(\mathbf{u}, r | \mathbf{u}_0, r_0) =$$
$$\left(- \sum_{i=1}^n \frac{\partial}{\partial u_i} D_i^{(1)} + \sum_{i,j=1}^n \frac{\partial^2}{\partial u_i \partial u_j} D_{ij}^{(2)} \right) p(\mathbf{u}, r | \mathbf{u}_0, r_0)$$

$$-\frac{\partial}{\partial r} \langle u^m v^n \rangle =$$
$$+ m \langle u^{m-1} v^n D_1 \rangle + n \langle u^m v^{n-1} D_2 \rangle$$
$$+ \frac{m(m-1)}{2} \langle u^{m-2} v^n D_{11} \rangle + \frac{n(n-1)}{2} \langle u^m v^{n-2} D_{22} \rangle$$
$$+ mn \langle u^{m-1} v^{n-1} D_{12} \rangle.$$

turbulence: long/transversal -3-


$$-r \frac{\partial}{\partial r} p(\mathbf{u}, r | \mathbf{u}_0, r_0) = \left(- \sum_{i=1}^n \frac{\partial}{\partial u_i} D_i^{(1)} + \sum_{i,j=1}^n \frac{\partial^2}{\partial u_i \partial u_j} D_{ij}^{(2)} \right) p(\mathbf{u}, r | \mathbf{u}_0, r_0)$$

$$- \frac{\partial \langle u^2 \rangle}{\partial r} = (2d_1^u + d_{11}^{uu}) \langle u^2 \rangle + d_{11} + d_{11}^{vv} \langle v^2 \rangle$$

$$- \frac{\partial \langle u^3 \rangle}{\partial r} = (3d_1^u + 3d_{11}^{uu}) \langle u^3 \rangle + 3d_{11}^u \langle u^2 \rangle + 3d_{11}^{vv} \langle uv^2 \rangle$$

Compare with the Karman/Horwarth equations:

$$- \frac{\partial \langle u^2 \rangle}{\partial r} = -2 \frac{\langle u^2 \rangle}{r} + 2 \frac{\langle v^2 \rangle}{r} \quad \text{First Kármán-}$$

$$- \frac{\partial \langle u^3 \rangle}{\partial r} = \frac{\langle u^3 \rangle}{r} - \frac{6 \langle uv^2 \rangle}{r} \quad \text{Second Kármán-}$$

turbulence: long/transversal -4-

rescaling symmetry: $r \Rightarrow 3r/2$

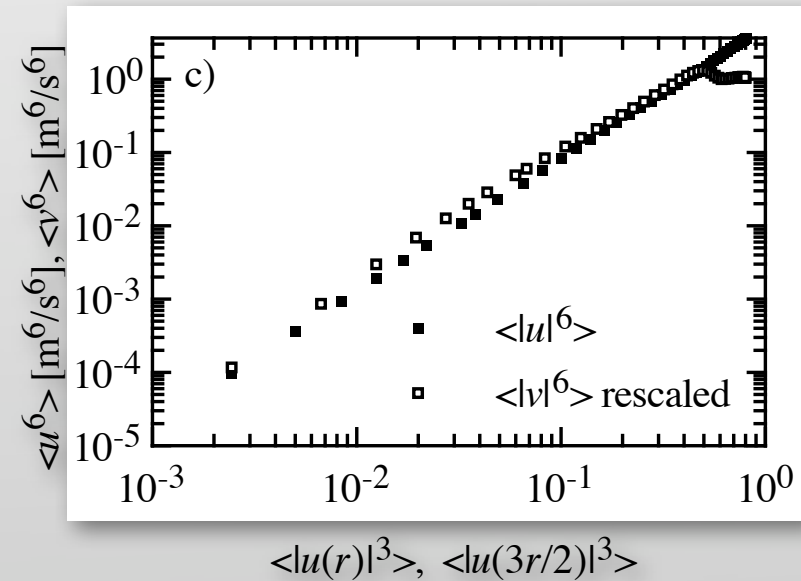
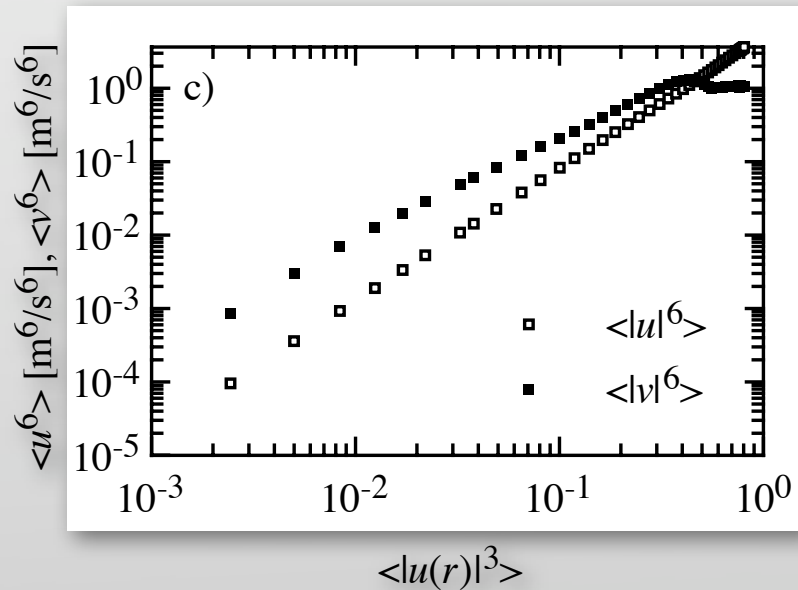
turbulence: long/transversal -4-

rescaling symmetry: $r \Rightarrow 3r/2$

$$\langle |v(r)|^n \rangle \propto \langle |u(r)|^3 \rangle^{\xi_n^t}$$

new ESST :

$$\langle |v(r)|^n \rangle \propto \langle |u(3r/2)|^3 \rangle^{\xi_n}$$



turbulence: long/transversal -5-

Assume • pure scaling: $\langle u^n(r) \rangle = c_l^n r^{\xi_l^n}$ and $\langle v^n(r) \rangle = c_t^n r^{\xi_t^n}$

• and same intermittency $d_{11}^{uu}(r) \equiv d_{22}^{vv}(r)$

$$\implies \langle v^n(r) \rangle = \langle u^n(\frac{3}{2}r) \rangle = c_t^n r^{\xi_t^n} = c_l^n (\frac{3}{2}r)^{\xi_l^n} \quad (\implies \xi_l^n = \xi_t^n)$$

$$n = 2 \quad c_t^2/c_l^2 \approx 1.33 \quad 2\% \text{ deviation from literature}$$

$$n = 4 \quad c_t^4/c_l^4 \approx 1.72 \quad 3\% \text{ deviation from literature}$$

turbulence: long/transversal -4-

rescaling symmetry: $r \Rightarrow 3r/2$

$$\langle |v(r)|^n \rangle \propto \langle |u(r)|^3 \rangle \xi_n^t$$

$$\langle |v(r)|^n \rangle \propto \langle |u(3r/2)|^3 \rangle \xi_n$$

consistent with Karman equation:

$$-r \frac{\partial}{\partial r} \langle u_r^2 \rangle = 2 \langle u_r^2 \rangle - 2 \langle v_r^2 \rangle$$

or

$$\langle v_r^2 \rangle = \langle u_r^2 \rangle + \frac{r}{2} \frac{\partial}{\partial r} \langle u_r^2 \rangle$$

taken as Taylor series

$$\langle v_r^2 \rangle \approx \langle u_{3/2r}^2 \rangle$$

turbulence: long/transversal -6-

rescaling symmetry: $r \Rightarrow 3r/2$

$$\langle |v(r)|^n \rangle \propto \langle |u(r)|^3 \rangle \xi_n^t$$

$$\langle |v(r)|^n \rangle \propto \langle |u(3r/2)|^3 \rangle \xi_n$$

striking result

$$\langle u_{3r/2}^n \rangle \propto (3r/2)^{\xi_n} \propto r^{\xi_n}$$

- this is only possible if the scaling laws does not hold

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- universal turbulence?
 - Reynolds dependencies

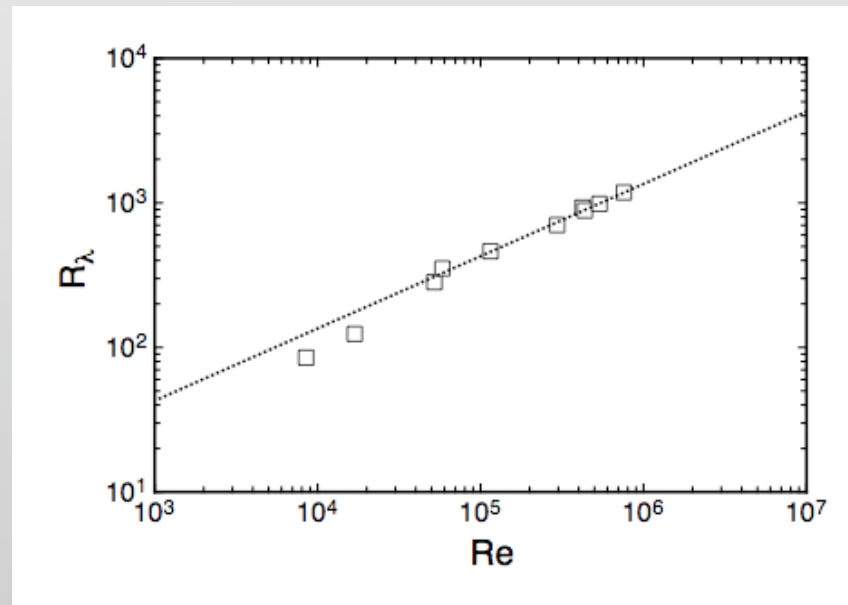
○ universality of turbulence: Re dependence?

$$D^{(1)}(u_r, r) = d_1^u(r) u_r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r) u_r + d_2^{uu}(r) u_r^2$$

Phys. Rev. Lett. **89**, (2002)

experimental data sets



Re dependence -2-

○ universality of turbulence: Re dependence?

$$D^{(1)}(u_r, r) = d_1^u(r) u_r$$

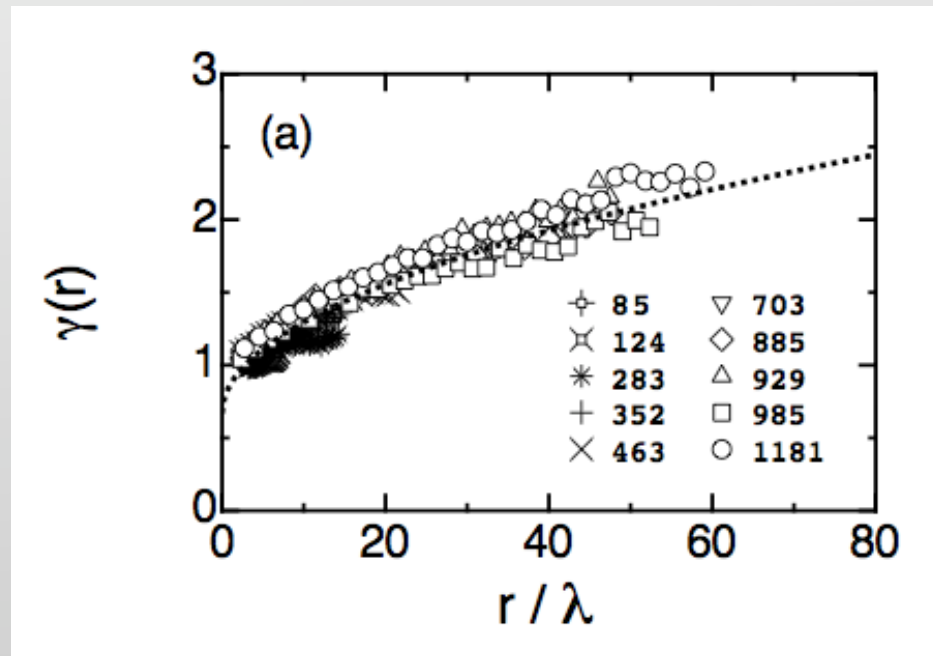
$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r) u_r + d_2^{uu}(r) u_r^2$$

Phys. Rev. Lett. **89**, (2002)

$$d_1^u(r) = -\frac{1}{r} \gamma(r)$$

$$\gamma(r) = \frac{2}{3} + c \sqrt{\frac{r}{\lambda}}$$

$D^{(1)}$ universal



Re dependence -3-

○ universality of turbulence: **Re dependence?**

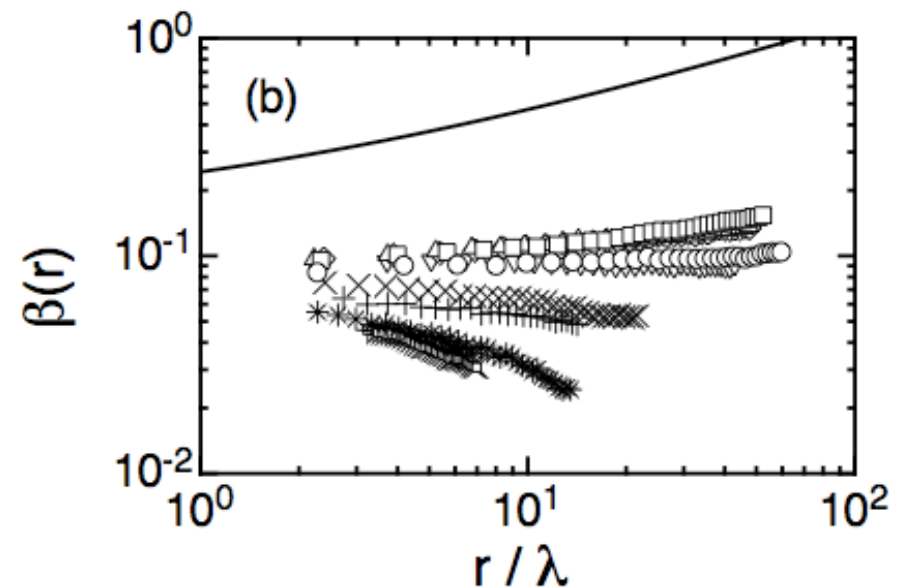
$$D^{(1)}(u_r, r) = d_1^u(r)u_r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r)u_r + d_2^{uu}(r)u_r^2$$

Phys. Rev. Lett. **89**, (2002)

$$d_2^{uu}(r) = f(Re)$$

$$d_2(r) \propto Re^{-3/8} / r$$



Re dependence -3-

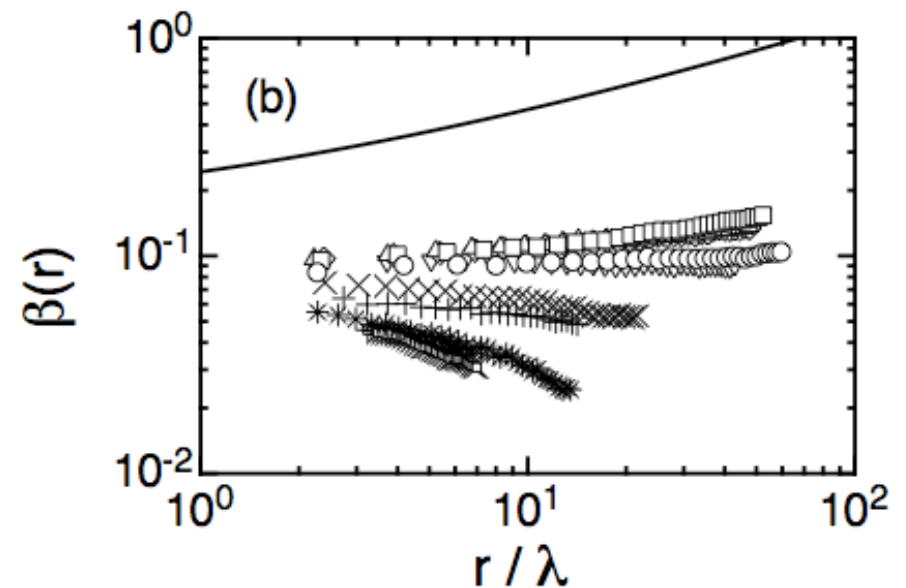
○ universality of turbulence: **Re dependence?**

$$\begin{aligned} D^{(1)}(u_r, r) &= d_1^u(r) u_r \\ D^{(2)}(u_r, r, Re) &= d_2(r, Re) + d_2^u(r, Re) u_r + d_2^{uu}(r, Re) u_r^2 \end{aligned}$$

Phys. Rev. Lett. **89**, (2002)

$$d_2^{uu}(r) = f(Re)$$

$$d_2(r) \propto Re^{-3/8} / r$$



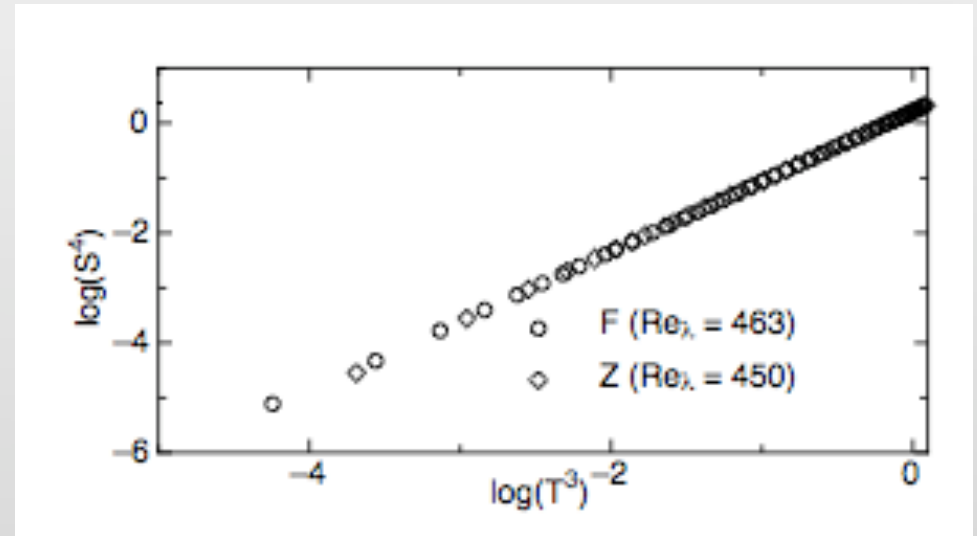
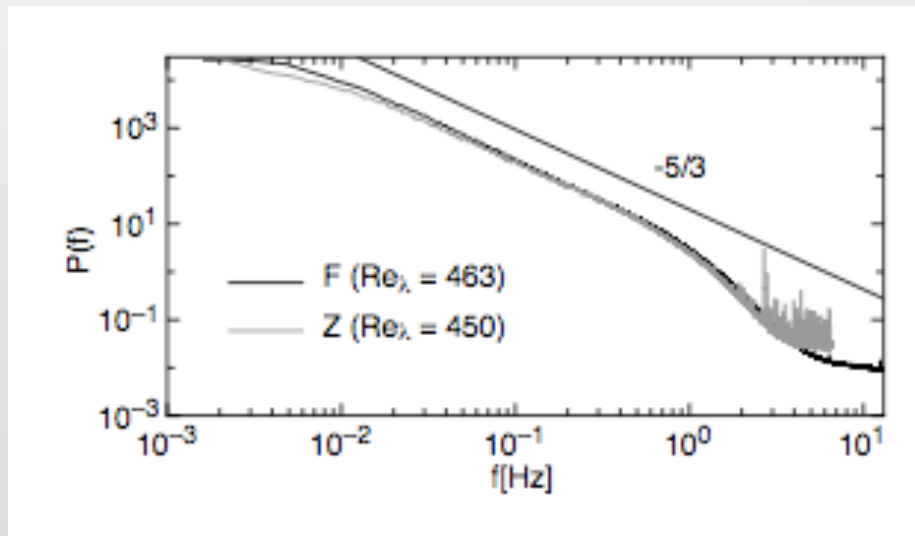
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- universal turbulence?
 - Reynolds dependencies
 - flow dependencies

universality of turbulence -2-

○ different **type of flows**

wake behind cylinder and free jet



Laupichler thesis. OI (2007)

universality of turbulence -2-

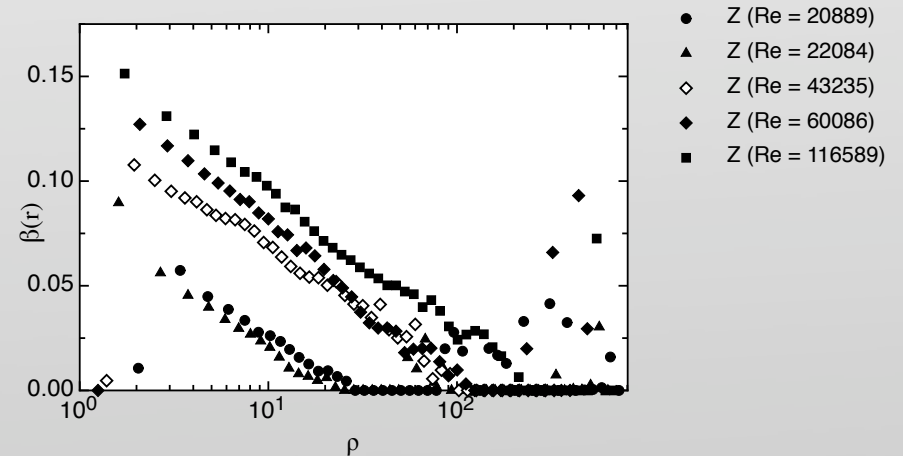
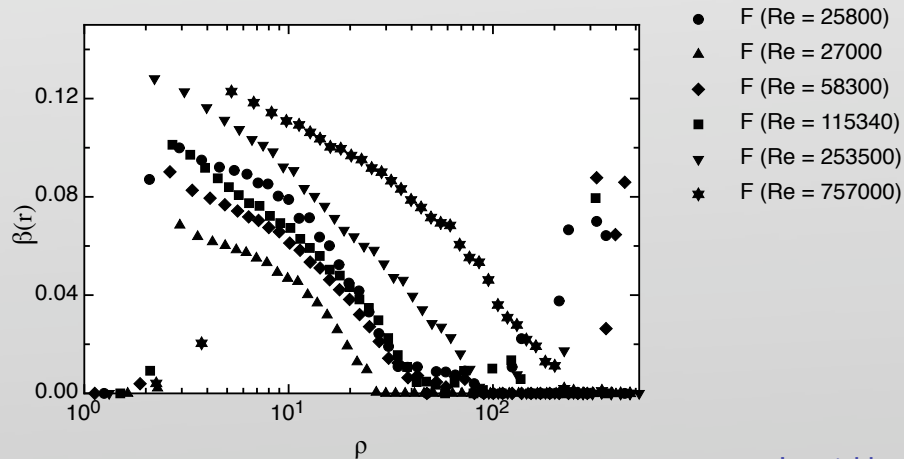
different type of flows

wake behind cylinder and free jet

$$-\frac{\partial}{\partial r}p(u_r, r|u_0 r_0) = \left[-\frac{\partial}{\partial u_r}D^{(1)}(u_r, r) + \frac{\partial^2}{\partial u_r^2}D^{(2)}(u_r, r)\right]p(u_r, r|u_0, r_0)$$

$$D^{(1)}(u_r, r) = d_1^u(r)u_r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r)u_r + d_2^{uu}(r)u_r^2$$

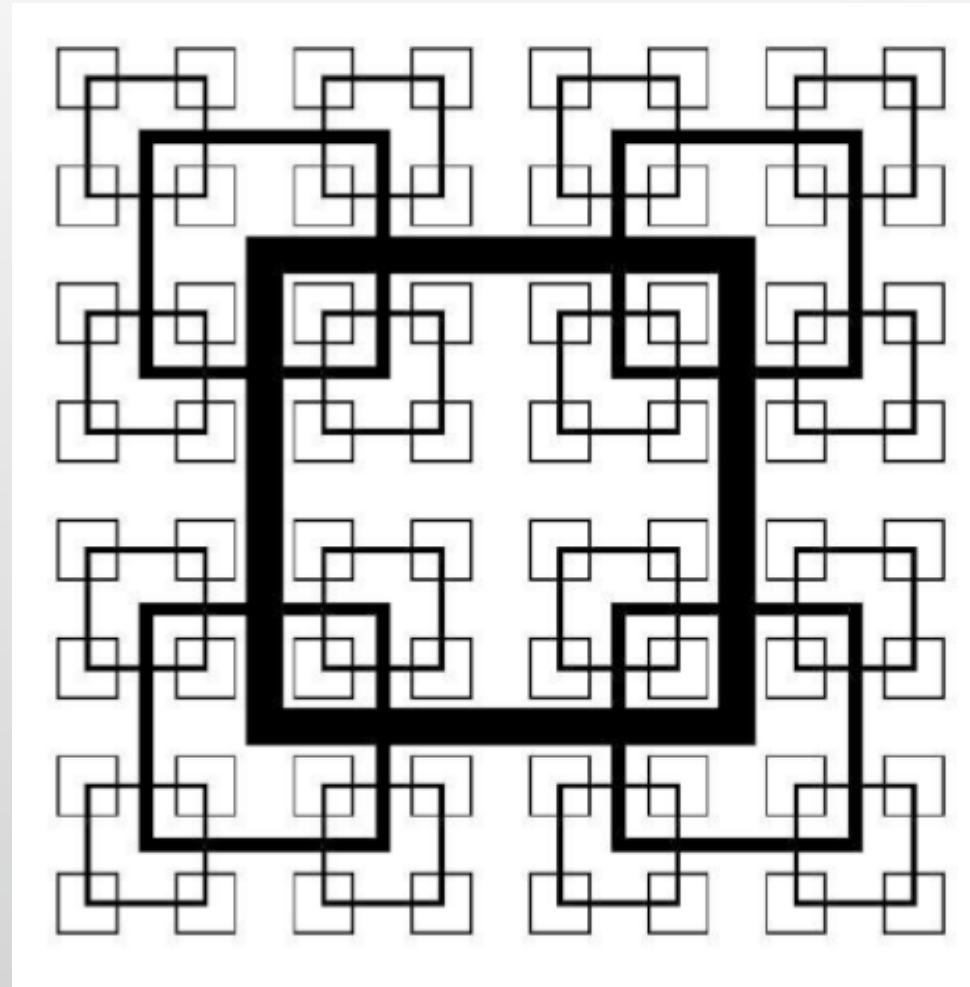


Laupichler thesis. OI (2007)

content

- comments on
 - universality is based on what?
- scaling properties
- characterization without scaling assumption
 - multi-scale statistics - additive noise
 - energy cascade and energy conservation
 - longitudinal-transversal structure functions
- universal turbulence?
 - Reynolds dependencies
 - flow dependencies

fractal grid turbulence - Christos Vassilicos

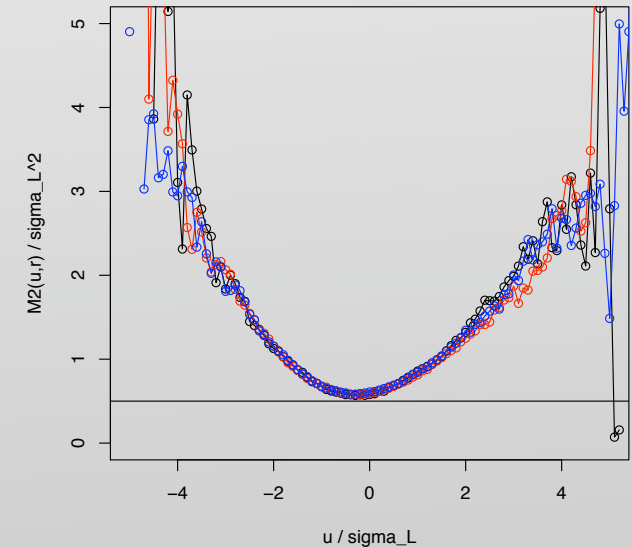
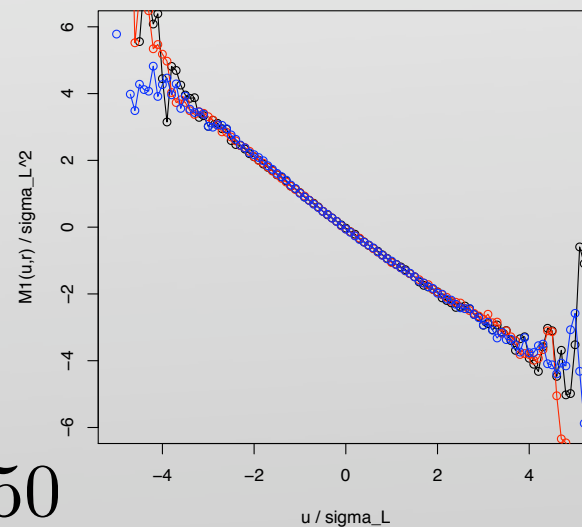
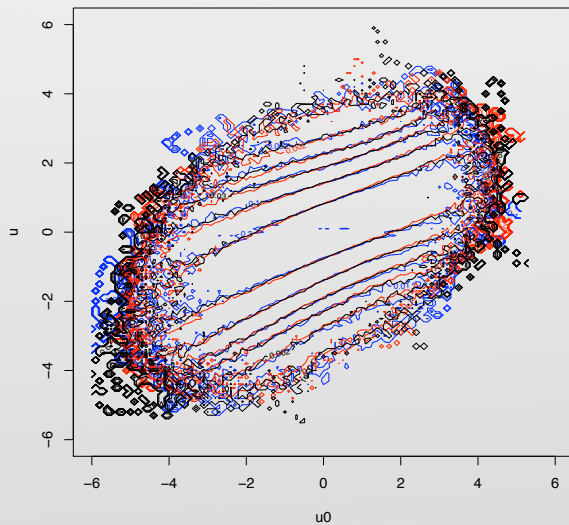


fractal grid turbulence - Christos Vassilicos

- properties do not change with Re

at L and L/2:

- conditional pdf
- Drift $D^{(1)}$ and diffusion $D^{(2)}$ do not depend on Re



$$R_\lambda = 150 \dots 750$$

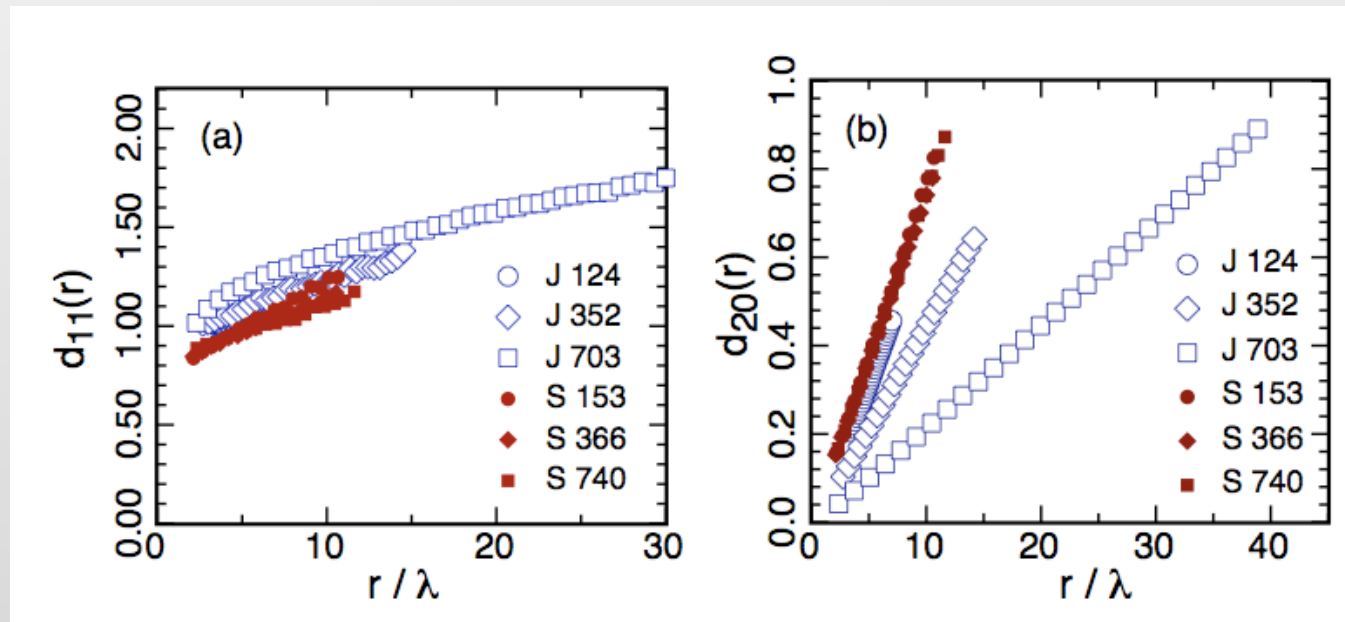
Phys. Rev. Lett. (2010) in press

fractal grid turbulence - Christos Vassilicos

○ Re dependency

$$D^{(1)}(u_r, r) = d_1^u(r) u_r$$

$$D^{(2)}(u_r, r) = d_2(r) + d_2^u(r) u_r + d_2^{uu}(r) u_r^2$$



=> Exp: cascade process depends on Re

Phys. Rev. Lett. (2010) in press

summary

- universality is based on what?
- scaling properties **one scale property**
- characterization with **multi-scale** statistics
 - additive noise
 - energy conservation
 - longitudinal-transversal structure functions
- universal turbulence?
 - Reynolds dependencies
 - flow dependencies
 - **new class of fractal generated turbulence**

**n-scale; n-point statistics
=> deeper insight into the nature of turbulence**

are there different classes of turbulence?